

# **Design of Microwave Resonant Cavity Transducer**

Development of Sensor Performance Model of Microwave Cavity Flow Meter for Advanced Reactor High Temperature Fluids

**Nuclear Science and Engineering Division** 

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## **Table of Contents**

Table of Contents	1
List of Figures	2
Abstract	3
1. Introduction	4
2. Sensor Model	6
2.1. Microwave resonant frequency shift	6
2.2. Thermophysical material properties	7
3. Computer Simulations of Sensor Performance	9
3.1. Selection of transducer geometric parameters	9
3.2. Computer simulations of microwave resonant frequency shift	10
3.3. Coupling of microwave field into cylindrical resonator	
4. Preliminary Transducer Design	15
5. Conclusions	17
References	18

# **List of Figures**

Figure 1 - (a) Conceptual drawing of waveguide-coupled microwave cylindrical cavity with
flexible membrane. (b) Deflection of membrane due to dynamic fluid pressure changes cavity
volume, which shifts microwave resonant frequency
<b>Figure 2</b> – Visualization of $TM_{011}$ mode in cylindrical cavity with $R$ =0.8in occurring at 15.4GHz.
Flexible membrane with $d=10$ mil is at the top of the cylinder
<b>Figure 3</b> – (a) Conceptual Figure 3. Frequency shift of $TM_{011}$ mode resonant frequency of $R$ =0.4in
and $d=10$ mil stainelss steel cylinder due to liquid sodium flow-induced membrane deflections at
$500^{\circ}\text{C}$ . Frequency shifts are calculated with COMSOL RF module using curved membrane model
Equation (2), with COMSOL RF using flat membrane model with equivalent amplitude $w_{\text{avg}}$ =
$w_0/3$ , and analytically using Equation (10).
Figure 4 – Visualization of TE <sub>011</sub> mode in right cylindrical cavity with $R=0.4$ in occurring at $f_0 =$
22.2GHz. Flexible membrane with $d=10$ mil is at the top of the cylinder
Figure 5 – Frequency shift of $TE_{011}$ mode resonant frequency shift of $R$ =0.4in and $d$ =10mil
stainless steel cylinder due to liquid sodium flow-induced membrane deflections at 500°C.
Frequency shift is calculated with COMSOL RF module using curved membrane model 12
Figure 6 – Visualization of $TE_{012}$ mode in right cylindrical cavity with $R=0.4$ in occurring at $f_0 =$
26.5GHz. Flexible membrane with $d=10$ mil is at the top of the cylinder
Figure 7 – Frequency shift of $TE_{012}$ mode resonant frequency shift of $R=0.4$ in and $d=10$ mil
stainless steel cylinder due to liquid sodium flow-induced membrane deflections at 500°C.
Frequency shift is calculated with COMSOL RF module using curved membrane model 13
Figure 8 – Dependence of Q factor of $TE_{012}$ mode on the radius of a circular aperture in the side
wall of the cylindrical cavity
<b>Figure 9</b> – Side view and 3D visualization of the cylindrical cavity design
Figure 10 – Top view of the cavity and 3D visualization of waveguide coupling cavity side wall

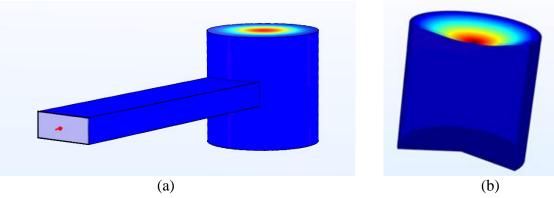
#### **Abstract**

High-temperature fluid reactors, such as molten salt cooled reactors (MSCR) and sodium fast reactors (SFR), are a promising advanced reactor option. Measurement of high-temperature fluid process variables, in particular the flow inside the pressure vessel, is a challenging task because of harsh environment, which includes high radiation, high temperature, and contact with highly corrosive coolant fluid. We are investigating a microwave cavity-based transducer for hightemperature fluid flow sensing. This sensor is a hollow metallic cylindrical cavity, which can be fabricated from stainless steel, and as such is expected to be resilient to radiation, high temperature and corrosive environment of MSCR and SFR. The principle of sensing consists of making one wall of the cylindrical cavity flexible enough so that dynamic pressure, which is proportional to fluid velocity, will cause membrane deflection. A cavity is characterized by its resonant frequencies. Membrane deflection causes cavity volume change, which leads to a shift in the resonant frequency. Feasibility of the flow sensor is evaluated with signal sensitivity using COMSOL computer simulations. A right cylinder geometry stainless steel cavity with dimeter of 0.8in was investigated. We choose membrane thickness of 10mil, so that corrosion anticipated to proceed at the rate of 1mil/year in liquid sodium would affect no more than 10% of the membrane. Using the properties of liquid sodium fluid, and stainless-steel material property values at 500°C, we calculate frequency shift for a range of values of fluid velocity from 0.5m/s to 2m/s. Results of computer simulations indicate measurable sensitivity to flow for this cavity design. Following these simulations, we have developed a preliminary design for fabrication of a transducer operating in microwave K-band for proof-of-principle tests.

#### 1. Introduction

High-temperature fluid reactors, such as sodium fast reactors (SFR) and molten salt cooled reactors (MSCR) are a promising advanced reactor option with highly efficient thermal energy conversion cycle [1-3]. Streamlining commercialization of advanced reactors involves development of new coolant sensing technologies for enhancing performance efficiency [4,5]. Measurement of high-temperature fluid process variables, in particular the flow inside the pressure vessel, is a challenging task because of harsh environments of advanced reactors, which includes radiation, high temperature, and contact with highly corrosive coolant fluid. Liquid metal flow sensors are typically of magnetic, Venturi and ultrasonic types [4,5], while molten salt flow sensors are ultrasonic [6]. Magnetic flow meters, which consist of high temperature permanent magnets (PM) and flux measurement coil (saddle coil), measure the rate of conducting flux passing through coil cross-section. The challenges of PM include heavy weight, temperature dependent output, and drifting with time. Difficulties with coil-based flow meters include large size of the DC power supply, temperature dependent output, nonlinearity at larger scales, and dependence on flow profile. Challenges with Venturi flow meters include long installation length, dependence on flow profile, and poor precision at low flow velocity. Ultrasonic flow meters are based on detecting either the time of flight or the Doppler frequency shift of acoustic waves in pitch-and-catch or in transmission mode. Ultrasonic transducers can be placed outside the pressure vessel. However, ultrasonic sensing requires a direct line-of-sight between the transducers. In addition, ultrasonic flowmeters are sensitive to flow profile in single path configuration, and thus might involve a multipath arrangement and thus complicated signal processing.

We are investigating a microwave cavity-based transducer for high-temperature fluid flow sensing. This sensor is a hollow metallic cavity, which can be fabricated from stainless steel, and as such is expected to be resilient to radiation, high temperature and corrosive environment of SFR and MSCR. A viable geometry of the sensor is that of a small cylindrical resonator [7,8]. A schematic drawing of the flow sensor is shown in Figure 1. The principle of sensing consists of making one wall of the cavity flexible enough so that dynamic pressure, which is proportional to fluid velocity, will cause membrane deflection. A cavity is characterized by its resonant frequencies, which occur due to constructive interferences of microwave field inside the cavity. Membrane deflection causes cavity volume change, and thus a shift in the resonant frequency. Using signal readout from microwave frequency shift in a hollow cavity is advantageous for applications in high temperature and high radiation environment, because no electronic components are placed inside the transducer. Since the microwave cavity sensing is based on fluidstructure interaction, as opposed to fluid electrical conductivity, the proposed sensor is equally applicable to liquid sodium and molten salt flow sensing. Energy coupling to and from the sensor will be achieved through a microwave waveguide, which will be an integral part of the insertion probe. A waveguide is a rigid narrow hollow metallic tube resilient to high temperature and highradiation environment. In the drawing in Figure 1, microwave field is coupled into the cylindrical resonant cavity through a side wall. This is typically accomplished by making a subwavelengthsize hole on the side of the cavity.



**Figure 1** – (a) Conceptual drawing of waveguide-coupled microwave cylindrical cavity with flexible membrane. (b) Deflection of membrane due to dynamic fluid pressure changes cavity volume, which shifts microwave resonant frequency.

The work described in this report consists of preliminary investigation of transducer design parameters, which is performed through analytic derivations and computer simulations with COMSOL software package. This includes a study of signal sensitivity to displacement of the membrane due to dynamic fluid pressure, and dependence of sensitivity on transducer geometrical parameters. Coupling of microwave field into the cavity is briefly discussed as well. Once the size of cavity and optimal resonant mode are identified, a cavity design is proposed for fabricating a cavity prototype for proof-of-principle tests.

#### 2. Sensor Model

#### 2.1. Microwave resonant frequency shift

The model of the sensor combines elements of fluid and solid mechanics, and electrodynamics. Dynamic pressure due to incompressible fluid flow normal to the membrane is given by Bernoulli equation:

$$P = \frac{1}{2}\rho v^2 \tag{1}$$

Where  $\rho$  and v are fluid density flow velocity, respectively. For uniform pressure, deflection of circular membrane of radius R as a function of radial position r can be calculated using Timoshenko's model for radially constrained circular plate [9]

$$w(r) = \frac{PR^4}{64D} \left( 1 - \frac{r^2}{R^2} \right)^2 = w_0 \left( 1 - \frac{r^2}{R^2} \right)^2$$
 (2)

Where D is material fexural rigidity given as

$$D = \frac{Ed^3}{12\left(1 - \nu^2\right)} \tag{3}$$

Here d is the plate thickness, E and v are the plate material Young's modulus and Poisson ratio, respectively. Maximum plate deflection occurs at the center

$$w(0) = w_0 = \frac{PR^4}{64D} = \rho v^2 \frac{3(1 - v^2)}{32E} \frac{R^4}{d^3}$$
 (4)

and the average plate deflection is given as

$$w_{avg} = \frac{1}{\pi R^2} \int_{0}^{2\pi} \int_{0}^{R} w(r) r dr d\phi = \frac{w_0}{3}$$
 (5)

Resonant frequencies of transverse electric  $TE_{nml}$  modes of a cylindrical cavity are characterized by integers n, m, l [7,8]

$$f_{nml}^{TE} = \frac{c}{2\pi\sqrt{\mu_r \varepsilon_r}} \sqrt{\left(\frac{X'_{nm}}{R}\right)^2 + \left(\frac{l\pi}{L}\right)^2}$$
 (6)

in which  $\mu_r$  and  $\varepsilon_r$  are permeability and permittivity of the medium, and  $X'_{nm}$  is the  $n^{th}$  root of the derivative of the  $m^{th}$  order Bessel function. The TM modes are described by a similar equation, with the difference that  $X'_{nm}$  is replaced with  $X_{nm}$ , the  $n^{th}$  root of the  $m^{th}$  order Bessel function. To prevent high-temperature fluid leaks into the sensor the cavity is pressurized with a buffer gas (e.g. Nitrogen or Helium). For gas-filled cavity,  $\mu_r=1$  and  $\varepsilon_r=1$ . The quality factor Q of a cavity is defined as

$$Q = \frac{f_0}{\Delta f_{\text{FWHM}}} \tag{7}$$

where  $f_0$  is the resonant frequency, and  $\Delta f_{FWHM}$  is the full width at half maximum spectral linewidth of the cavity. The Q of a gas-filled cavity is related to resisitve losses at the walls, which are determined by the material electrical conductivity. When optimizing Q with respect to cavity geometry, the maximum Q value for low order modes such as  $TE_{01p}$ , is achieved when L=2R. The cavity resonant frequency can be characterized by a single parameter

$$f_{nmp}^{TE} = \frac{c}{2\pi L} \sqrt{(2X'_{nm})^2 + (l\pi)^2} = \frac{c}{2L} g^{TE}(n, m, l)$$
 (8)

and simimarly for TM mode frequencies, where  $g^{TE}(n,m,l)$  and  $g^{TM}(n,m,l)$  are non-dimensional function of the mode number integers n,m,l. Typical value of Q for the low order mode of the cylindrical cavity is approximately 20,000.

For small changes in the cavity hight due to membrane displacement of magnitude  $\Delta L$ , frequency shift is given, to first order in  $\Delta L$ , as

$$\Delta f = \frac{df}{dL} \Delta L = -\frac{c}{2L^2} \Delta L g(n, m, l) \tag{9}$$

For  $\Delta L = -w_{\text{avg}} = -w_0/3$ , and using L = 2R,

$$\Delta f = g(n, m, l) \frac{1 - v^2}{256E} \frac{cR^2}{d^3} \rho v^2$$
 (10)

With respect to geometrical parmeters of the cavity, detection sensitivity increases quadratically with R, and inverse cubic with d. With respect to fluid parameters, detection sensitivity increases linearly with density and quadratically with velocity.

### 2.2. Thermophysical material properties

The properties of materials and fluids of variables which enter Equation (10) are functions of temperature. For stainless steel 316, Young's modulus E, shear modulus G, and Poisson ratio v are given as [10]

$$E = 2.137 \times 10^5 - 102.74T \tag{11a}$$

$$G = 8.964 \times 10^4 - 53.78T \tag{11b}$$

$$v = E/2G - 1 \tag{11c}$$

Where *E* and *G* are given in MPa, *T* is given in  $^{\circ}$ C, and the correlations are valid for  $0 < T < 800^{\circ}$ C. Temperature-dependent values of liquid sodium density  $\rho$  in units of kg/m<sup>3</sup> is given by the correlation [11]

$$\rho = \rho_c + f(1 - T/T_c) + g(1 - T/T_c)^h \tag{12}$$

Valid for 371K<T<2503.7K, where  $\rho_c$ =219, f=275.32, g=511.58, h=0.5,  $T_c$ =2503.7K.

For molten salt, density decreases linearly with temperature. For example, for FLiBe (LiF-BeF<sub>2</sub>), density correlation in units of kg/m<sup>3</sup> is given by [12]

$$\rho = 2415.6 - 0.49072T \tag{13}$$

Where *T* is in units of K. The correlation is valid for 732.2K<*T*<4498.8K.

### 3. Computer Simulations of Sensor Performance

#### 3.1. Selection of transducer geometric parameters

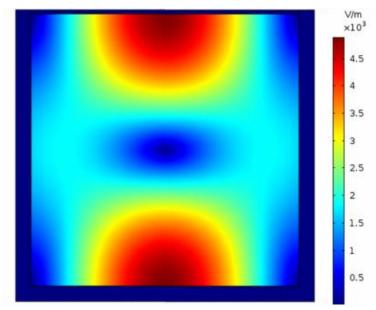
Thickness of the membrane is a critical parameter that affects sensor sensitivity. The value of thickness of the membrane is selected as a tradeoff between sensitivity and structural integrity of the sensor. One of the possible pathways for sensor failure is due to corrosion of the membrane. Prior studies indicate that stainless steel 316 is highly resilient to corrosion in liquid sodium. The data is usually based on converting the weight change of thin samples to loss rates in mils per year, and thus might not account for local inter-granular penetration damage which could affect the mechanical properties. The data indicates that corrosion of stainless steel in liquid sodium depends on the concentration of oxide impurities, sodium flow velocity, and temperature. One study showed that, when the maximum temperature of the liquid sodium is 600°C, the oxide level is less than 10 ppm, and the velocity is 2m/s, the corrosion rate at 2500 hours is approximately 2.3µm/year [13]. If the velocity is increased to 6m/s, the corrosion rate is approximately 2.8µm/ year, indicating that the velocity has only a slight effect on the corrosion rate. However, if the velocity is held at 6m/s, and the maximum temperature is increased to 720°C, the corrosion rate is increased roughly by one order of magnitude to approximately 30.2µ/year. In the case of molten salt, 2µm/year corrosion rate of stainless steel 316 was observed in purified FLiBe (2LiF-BeF<sub>2</sub>) at 650°C in a static cell (no flow) [14,15]. However, another study of SS316 in FLiBe estimated the corrosion rate at 700°C estimated the corrosion rate to be between 17.1µm/year and 31.2µm/year. In the liquid sodium environment of the Mechanisms Engineering Test Loop (METL) [16], the operating conditions are maximum temperature of 650°C, liquid sodium velocity in pipes under 1m/s, and oxide concentration of <5ppm. The corrosion rate of SS316 is expected to be approximately 0.1mil/year. Thus, in a 10year life span, the sensor membrane, on average, could corrode by 1mil. We choose d=10mil (254µm), so that corrosion would affect no more than 10% of the membrane. This number corresponds to a commonly available stainless steel stock sheet.

Another parameter that determines sensor sensitivity is the radius of the membrane. Besides the quadractic dependence on the radius in Equation (10), for a given mode, the size of the resonator determines the values of the resonant frequency. Fluid flow sensing is accomplished by measuring the shift of the resonant frequency caused by membrane displacement. The measurement procedure would involve scanning over the frequency span of  $\Delta f_{FWHM}$  centered at  $f_0$ , with frequency resolution bandwidth  $f_{RBW}$  of the spectrum analyzer readout. The trace sweep time  $t_{sweep}$  would contribute to the flow meter time constant. Note that for a given value  $f_{RBW}$  and  $t_{sweep}$  are inversely related. For example, for of L=2R=0.8in (2cm) gives  $f_0\sim20$ GHz. Using Q = 20,000, we obtain and  $\Delta f_{FWHM}\sim1$ MHz. For such parameters, a modern state-of-the-art power network analyzer (PNA) can perform the frequency sweep with  $f_{RBW}=3$ KHz and  $t_{sweep}=62$ ms. The condition for dectability is  $\Delta f > f_{RBW}$ .

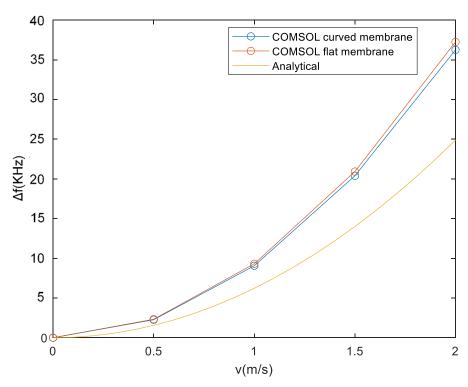
#### 3.2. Computer simulations of microwave resonant frequency shift

To confirm validity of the analytical model approximations, frequency shift in a microwave cavity was calculated with COMSOL RF module. The modes of cylindrical cavity in order of increasing frequency are  $TM_{010}$ ,  $TM_{011}$ ,  $TM_{012}$ ,  $TE_{011}$ ,  $TE_{012}$ , etc. The lowest frequency mode  $TM_{010}$  has no variation of the E-field with length, and therefore it is almost insensitive to length change. The subsequent modes, however, are sensitive to cavity volume change resulting from membrane displacement. COMSOL simulation results for a R=0.4in and d=10mil membrane cylindrical cavity in Figure 2 show spatial distribution of the electric fields in the cavity for  $TM_{011}$  mode, occurring at 15.4GHz. The flexible membrane is at the top of the cylinder.

Using stainless-steel and liquid sodium material property values at  $500^{\circ}$ C given by Equations (11) and (12), respectively, we estimate  $\Delta f$  for a range of values velocity values from 0.5m/s to 2m/s. Frequency shifts were calculated with COMSOL RF module for TM<sub>011</sub> mode. COMSOL simulations were performed for parabolic membrane deflection profile given by Equation (2), and for flat membrane displacement with amplitude  $w_{avg} = w_0/3$  according to Equation (5). For comparison, frequency shift was also calculated with the analytical model of Equation (10). The results are plotted in Figure 3. Notice that the COMSOL RF calculations for parabolic membrane deflection shape and for equivalent amplitude flat membrane displacement agree closely, and also agree reasonably well with frequency shifts calculated with the analytical model of Equation (10).



**Figure 2** – Visualization of TM<sub>011</sub> mode in cylindrical cavity with R=0.8in occurring at 15.4GHz. Flexible membrane with d=10mil is at the top of the cylinder

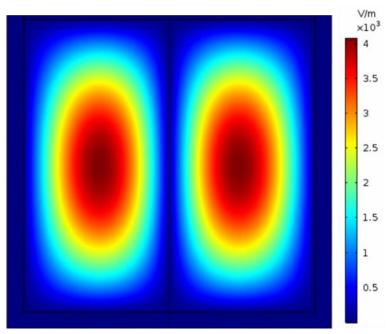


**Figure 3** – (a) Conceptual Figure 3. Frequency shift of TM<sub>011</sub> mode resonant frequency of R=0.4in and d=10mil stainelss steel cylinder due to liquid sodium flow-induced membrane deflections at 500°C. Frequency shifts are calculated with COMSOL RF module using curved membrane model Equation (2), with COMSOL RF using flat membrane model with equivalent amplitude  $w_{avg}$ =  $w_0/3$ , and analytically using Equation (10).

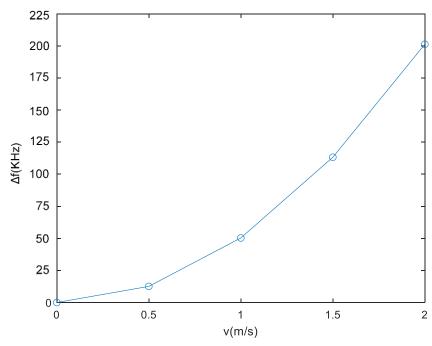
Higher frequency modes, such as  $TE_{011}$  and  $TE_{012}$  modes, are known to have higher Q values than TM modes. Thus, we expect  $TE_{011}$  and  $TE_{012}$  modes to have better sensitivity to membrane displacement. Signal sensitivity for these two modes has been investigated through computer simulations, with results shown in Figures 4,5,6, and 7. COMSOL simulation results for R=0.4in and d=10mil right cylindrical cavity in Figure 4 show spatial distribution of the electric fields in  $TE_{011}$  mode, occurring at 22.2GHz resonant frequency with Q = 17,900. Frequency shifts of the stainless steel cavity were calculated with COMSOL RF module for  $TE_{011}$  mode for a range of velocities (0.5m/s to 2m/s) of liquid sodium flow at 500°C. The membrane defection profile was modeled with Equation (2). Results of frequency shift computer simulations are shown in Figure 5. As expected, for the same fluid velocity, frequency shifts of  $TE_{011}$  mode are larger than those of the  $TM_{011}$  mode.

COMSOL simulation results for R=0.4in and d=10mil right cylindrical cavity in Figure 6 show spatial distribution of the electric fields in TE<sub>012</sub> mode, occurring at 26.5GHz resonant frequency with Q = 19,500. Frequency shifts for TE<sub>012</sub> mode for a range of velocities (0.5m/s to 2m/s) of liquid sodium flow at 500°C are displayed in Figure 7. Because of the higher Q value and spatial distribution profile, frequency shifts for TE<sub>012</sub> are higher than those of TE<sub>011</sub> mode. As discussed

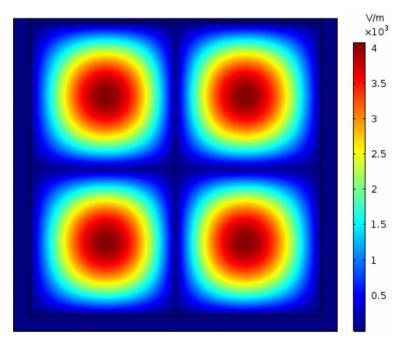
in Section 3.1, frequency shifts larger than 3KHz are, in principle, detectable with modern electronic devices.



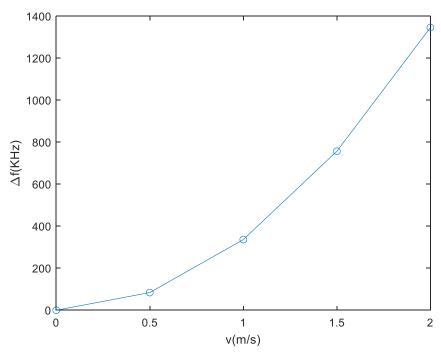
**Figure 4** – Visualization of TE<sub>011</sub> mode in right cylindrical cavity with R=0.4in occurring at  $f_0$  = 22.2GHz. Flexible membrane with d=10mil is at the top of the cylinder



**Figure 5** – Frequency shift of  $TE_{011}$  mode resonant frequency shift of R=0.4in and d=10mil stainless steel cylinder due to liquid sodium flow-induced membrane deflections at 500°C. Frequency shift is calculated with COMSOL RF module using curved membrane model.



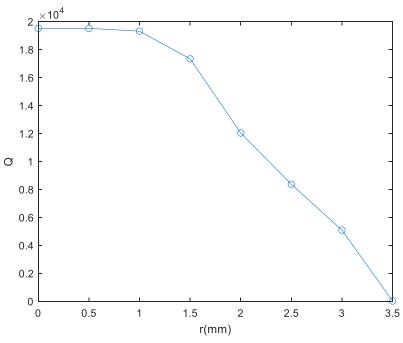
**Figure 6** – Visualization of  $TE_{012}$  mode in right cylindrical cavity with R=0.4in occurring at  $f_0$  = 26.5GHz. Flexible membrane with d=10mil is at the top of the cylinder



**Figure 7** – Frequency shift of  $TE_{012}$  mode resonant frequency shift of R=0.4in and d=10mil stainless steel cylinder due to liquid sodium flow-induced membrane deflections at 500°C. Frequency shift is calculated with COMSOL RF module using curved membrane model

#### 3.3. Coupling of microwave field into cylindrical resonator

An important consideration for transducer design is the strategy for coupling microwave field into the cavity. As discussed in Section 1, our envisioned approach consists of coupling microwave field into the cylindrical resonator through a subwavelength-size aperture in the side wall of the cylinder. Selecting the size of an aperture involves a trade-off between coupling sufficient microwave energy into the cavity and degrading Q factor of the cavity due to microwave leakage. Based on results of computer simulations in Section 2.2, best sensitivity to fluid velocity is expected for TE<sub>012</sub> resonant mode. Considering spatial distribution of TE<sub>012</sub>, which is visualized in Figure 6, there is an electric field node in the mid-plane of the cavity. Thus, we anticipate that Q-factor of TE<sub>012</sub> would display reasonable resilience to the size of the aperture. Quantitative dependence of Q factor on the aperture size for the TE<sub>012</sub> is investigated with COMSOL RF module computer simulations, the results of which are shown in Figure 8. Note that the free space wavelength of the  $TE_{012}$  is  $\lambda_0=1.13$ cm. One can observe in Figure 8 that the Q factor does not decrease significantly for r<1.5mm, or r<0.1 $\lambda_0$ . As r increases above 1.5mm value, Q factor of the cavity is rapidly decreasing. This suggests the upper bound value on the coupling aperture value for the transducer design. Note that excitation of TE<sub>012</sub> frequency in a cylindrical cavity can be accomplished with commercially available K-band (18-26.5GHz) metallic waveguides.



**Figure 8** – Dependence of Q factor of  $TE_{012}$  mode on the radius of a circular aperture in the side wall of the cylindrical cavity

### 4. Preliminary Transducer Design

Based on the results obtained from modeling and simulation of resonant transducer performance, preliminary design of the microwave cavity transducer was developed for fabrication in a machine shop. Because the resonant frequency of TE<sub>012</sub> resonant mode discussed in Section 3 is in the K-band, the dimensions of the cylindrical right cavity were slightly adjusted to match those of the standard W/UG-595/U WR-42 waveguide flange (0.875in x 0.875in square). Slight increase in cavity dimensions will result in slight decrease of the resonant frequency modes. The side view of the cavity is shown in Figure 9. All dimensions are in units of inches. In this design, the cavity is assembled with screws from a hollow cylinder, a thick cap at the bottom, and a thin diaphragm made from 10mil shim stock metal. The thin diaphragm is held in place with a clamping ring. All the holes in the drawing are for 4-40 screws, which are the same size screws used for Kband waveguide flange. In the design shown in Figure 6, microwave field is coupled into the cavity through a subwavelength aperture on the side of the cylinder. The diameter of the aperture is 3mm (0.116in), which is the largest diameter that does not significantly degrade the Q factor of TE<sub>012</sub> mode according to computer simulations in Section 3.3. The aperture is matched to the center of the waveguide. Coupling of waveguide to the cylindrical cavity is accomplished by chamfering the side of a cylindrical cavity to obtain a flat surface. 3D rendering of the cylindrical cavity in Figure 9 shows the chamfered side surface and the small aperture for coupling microwave field into cavity.

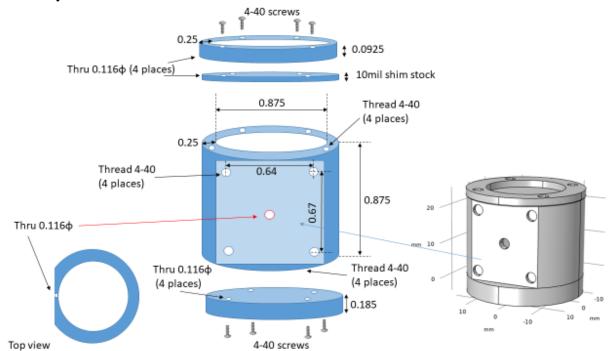


Figure 9 – Side view and 3D visualization of the cylindrical cavity design

Visualization of WR-42 microwave K-band waveguide attached to cylindrical cavity is shown in Figure 10. The top view of the cylinder without the diaphragm is also shown in Figure 10. In

addition to the diameter of the aperture, another parameter influencing efficiency of microwave coupling into cavity is the thickness of the aperture screen. Selecting aperture thickness involves a tradeoff between coupling efficiency and mechanic structural integrity of the resonator. In the design shown in Figure 10, thickness of the aperture is 0.0925in (2.3mm). Dependence of coupling efficiency on screen thickness will be further investigated through computer simulations and proof-of-principle tests with coupons with apertures of different thicknesses.

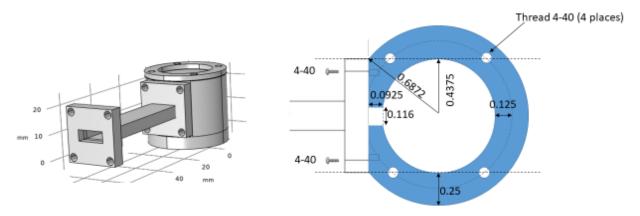


Figure 10 – Top view of the cavity and 3D visualization of waveguide coupling cavity side wall

#### 5. Conclusions

We have performed preliminary investigation of performance of a microwave cavity-based transducer for high-temperature fluid flow sensing using analytical derivations and computer simulations. This sensor is a hollow metallic cylindrical cavity, which can be fabricated from stainless steel, and as such is expected to be resilient to radiation, high temperature and corrosive environment of MSCR and SFR. The principle of sensing consists of making one wall of the cylindrical cavity flexible enough so that dynamic pressure, which is proportional to fluid velocity, will cause membrane deflection. A cavity is characterized by its resonant frequencies. Membrane deflection causes cavity volume change, which leads to a shift in the resonant frequency. Feasibility of the flow sensor is evaluated with signal sensitivity using COMSOL computer simulations. A right cylinder geometry stainless steel cavity with dimeter of 0.8in was investigated. We choose membrane thickness of 10mil, so that corrosion anticipated to proceed at the rate of 1mil/year would affect no more than 10% of the membrane. Using the properties of liquid sodium fluid, and stainless-steel material property values at 500°C, we calculate frequency shift for a range of values of fluid velocity from 0.5m/s to 2m/s. Results of computer simulations indicate measurable sensitivity to flow for this transducer design. Computer simulations indicated the best sensitivity to flow-induced membrane displacement will be achieved when the cavity is excited in TE<sub>012</sub> mode. Coupling of microwave filed into the cavity can be accomplished through a subwavelength size circular aperture on the cavity side wall. Computer simulations indicate that the Q factor of the cavity is not affected significantly for cavity radius r<0.1 $\lambda_0$  (free space wavelength of TE<sub>012</sub> mode. Following computer simulation results, transducer design was developed for fabrication of cavity prototype for proof-of-principle tests. For the design parameters identified in this study through computer simulation, the prototype cavity will operate in microwave K-band (18-26.5GHz). This cavity will be coupled with WR-42 K-band microwave waveguide. Future work will focus on sensor fabrication and tests to validate performance in a fluid.

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